# **Experimental methods**

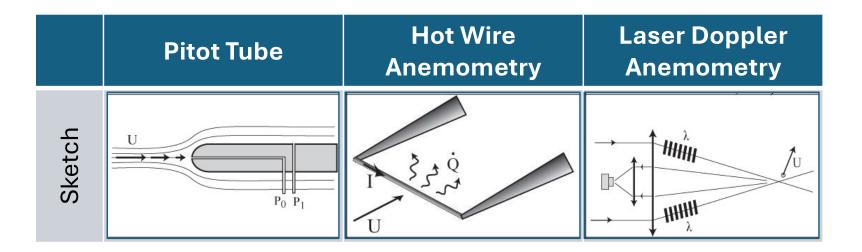
# One point techniques

**Romain MONCHAUX** 

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Institut Polytechnique de Paris

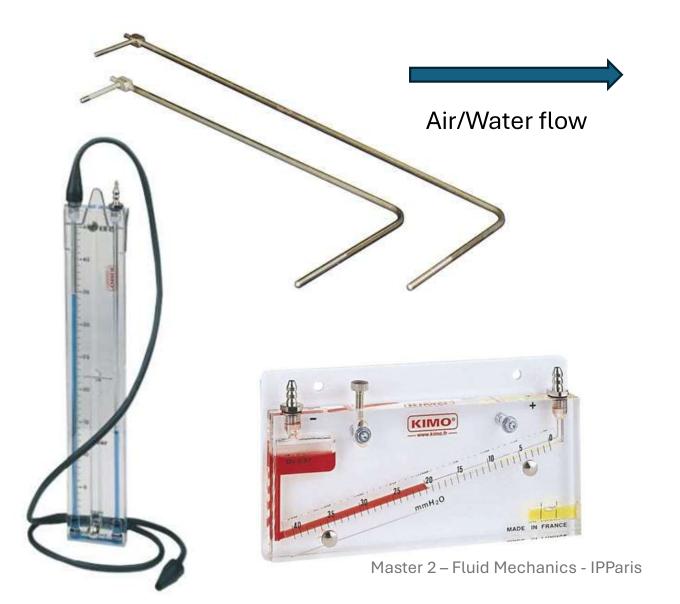
## One point techniques

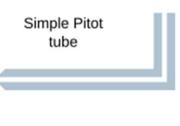


## Pitot tube

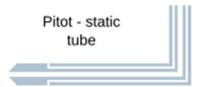
## **Principle:**

Stagnation pressure measurement Static pressure measurement Bernoulli theorem











## Pitot tube

#### Principle:

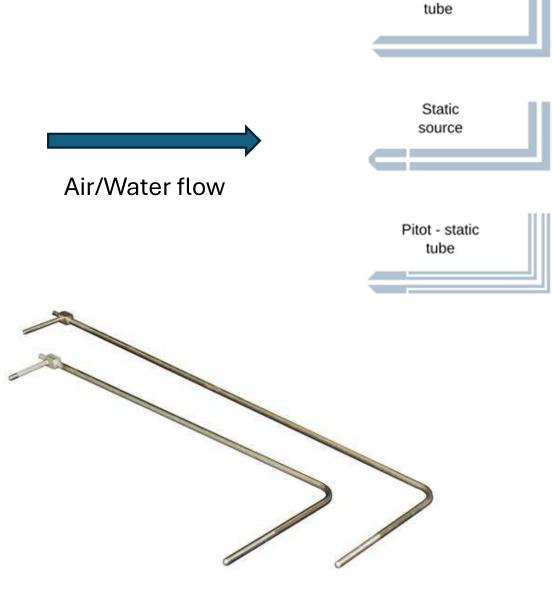
Stagnation pressure measurement Static pressure measurement Bernoulli theorem

#### Pros / cons:

Cheap
Easy to use
Suited for time averages
Intrusive
Poor time resolution
Poor spatial resolution

#### **Uses:**

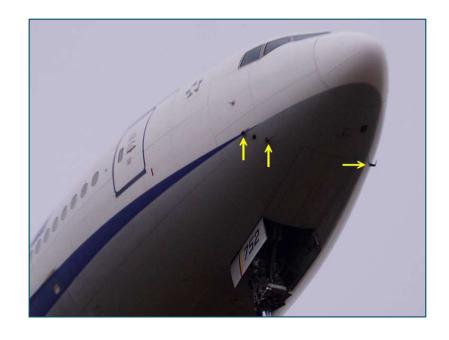
Monitoring Calibration In situ Large arrays



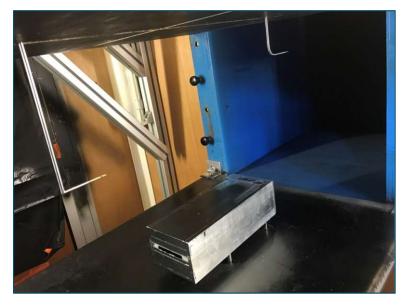
Simple Pitot

# Pitot tube



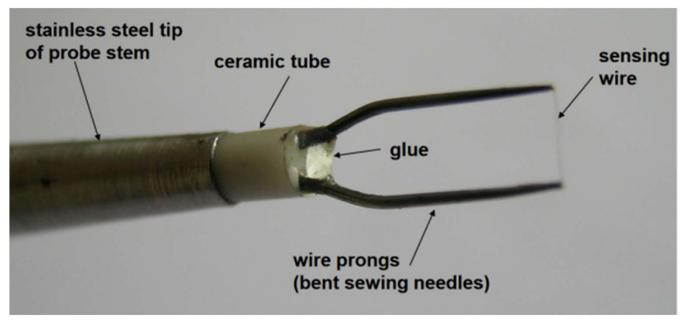




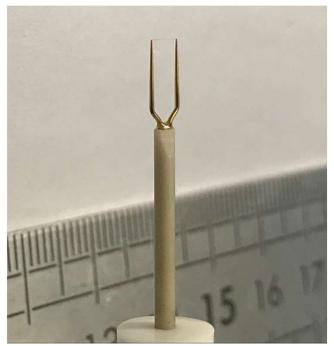


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#### How does it look like?



Platinium or tungsten thin wire (few microns thick, few mm long) welded to the prongs of the probe support



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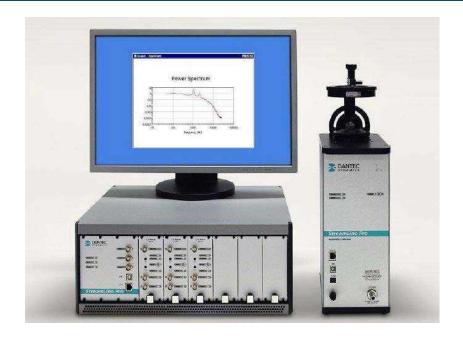
Illustrations: von Karman institute, ENSTA

#### How does it look like?



Pocket size model (mean velocity)

150 – 400 €



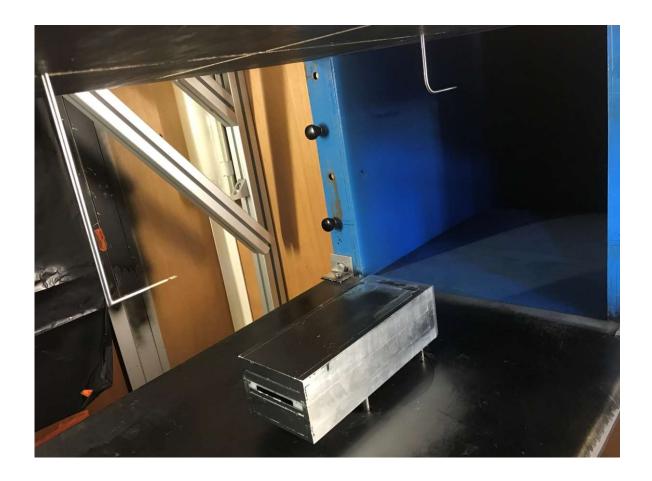
Lab model (fluctuations, 10 probes) 30-40 k€

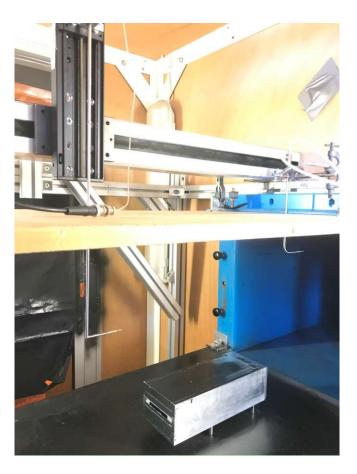


Mini CTA (fluctuations, 1 probe) 2-4 k€

Illustrations: DANTEC dynamics

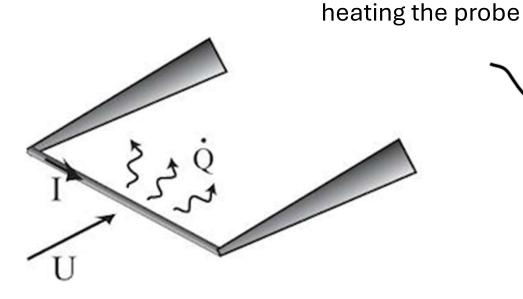
#### How does it look like?



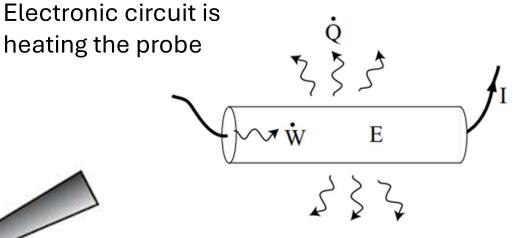


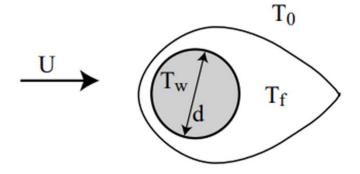
ENSTA wind tunnel – Ahmed body + hot wire

#### How does it work?



Flow is cooling the probe: « chill factor »





Thermal boundary layer



#### **Governing equations**

Energy stored in the wire:

$$\frac{dE}{dt} = \dot{W} - \dot{Q}$$

Electrical power  $\dot{W} = R_w I^2 > 0$ 

Joule effect

#### $\dot{Q} > 0$ Heat variation

Conduction to fluid

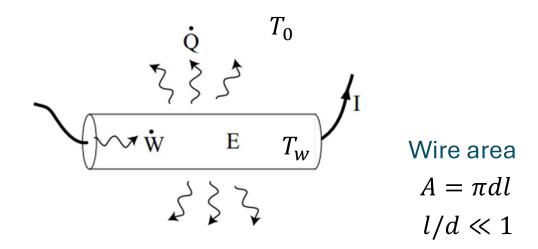
**Convection to fluid** 

Conduction to support

Radiation

## Heat flux

$$\overrightarrow{j_Q} = -k_f \overrightarrow{\nabla} \overrightarrow{T}$$



#### Nusselt number

 $Nu = hd/k_f$ 

Convection efficiency

$$Nu = I$$

$$\longrightarrow \dot{Q} = k_f l\pi \left( T_w - T_0 \right) Nu$$

 $\frac{\dot{Q}}{A} = h(T_w - T_0)$ 

#### **Governing equations**

Power budget 
$$\frac{dE}{dt} = R_w I^2 - \pi l k (T_w - T_0) N u$$

$$R_w I^2 = \pi l k (T_w - T_0) N u$$

## Link between Nu and the velocity U?

#### **Governing equations**

Heat transfer: Nu = f(Re, Pr, Ma, l/d ...)

Wire Reynolds number: 
$$Re_w = \frac{Ud}{v_f} \ll 1$$

Fluid Prandtl number: 
$$P_r = \frac{v_f}{\kappa_f} \simeq 1$$



Laminar flow

Thermal and viscous boundary layers similar

$$Nu = 1 + \sqrt{2\pi Re_w}$$

King's laws 
$$Nu = 0.42 P_r^{\frac{1}{5}} + 0.57 Re_w^{\frac{1}{2}} P_r^{\frac{1}{3}}$$

$$Nu=a_0+b_0\sqrt{Re_w}$$

#### **Governing equations**

$$R_w I^2 = \pi l k (T_w - T_0) (a_0 + b_0 R e_w^{\frac{1}{2}})$$

What do we measure?

- Tension:  $e = R_w I$
- No link with U a priori
- But ...

$$R_w = f(T_w)$$

$$R_{w} = R_{0}(1 + \beta(T_{w} - T_{0}))$$

$$\beta = \frac{1}{R_{w}} \frac{\partial R_{w}}{\partial T} \simeq cst$$

$$R_{w}I^{2}$$

$$R_{w}I^{2}$$

$$R_{w} - R_{0}$$

For Tu or Pt:

$$\beta \simeq 5 \times 10^{-3} K^{-1}$$

$$\frac{e^2}{R_W(R_W - R_0)} = a + b\sqrt{U}$$

#### **Two strategies**

$$\frac{e^2}{R_w(R_w - R_0)} = a + b\sqrt{U}$$

$$e = R_w I$$

King's law

Ohm's law

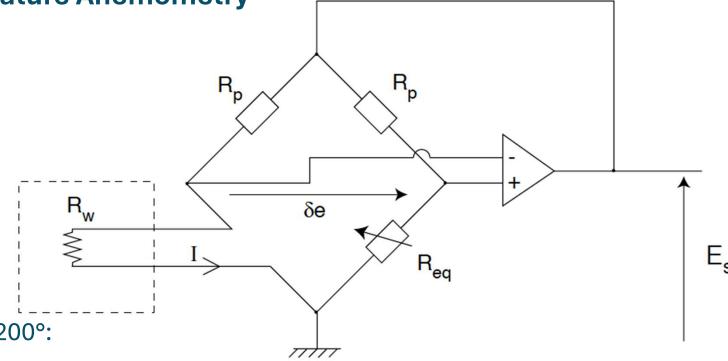
### **Constant current anemometry CCA**

- *I* is constant
- U through fluctuations of  $R_w$
- Obsolete

### **Constant temperature anemometry CTA**

- $R_w$  i.e.  $T_w$  is constant
- U through fluctuation of I i.e. e





Overheat 100° - 200°:

$$R_{eq} < R_w$$

Negative feedback:

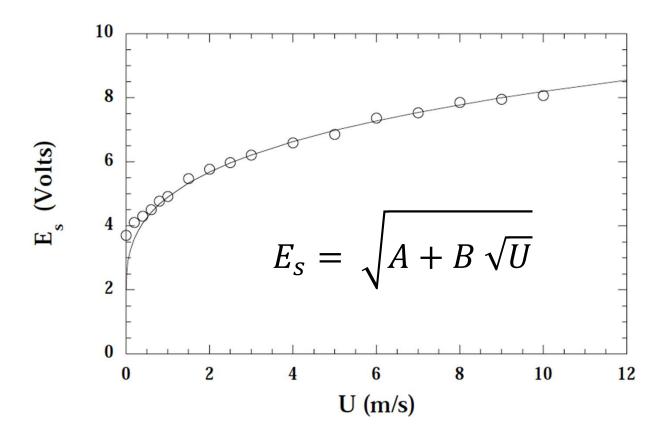
$$\delta e = I(R_w - R_0)$$

$$E_{\rm S} = G \, \delta e$$

$$E_S = \sqrt{R_W(R_W - R_0)(a + b\sqrt{U})}$$

$$E_s = \sqrt{A + B\sqrt{U}}$$

#### **Calibration**



#### **Modified King's law**

$$E_s^2 = A + BU^n$$

n between 0.4 and 0.6

### **Fitting strategy**

$$A = E_s^2(0)$$
  
 
$$n \log U + \log B = \log \left( E_s^{2(U)} - A \right)$$

#### **Empirical relations**

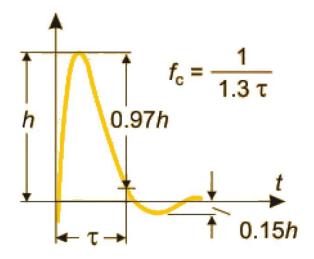
Collis & Williams (1959)'s empirical relation

$$n = 0.45$$
 for  $0.02 < Re_w < 44$   
 $n = 0.51$  for  $44 < Re_w < 140$ 

#### The wire is a cylinder in a cross flow:

- $Re_w < 44$ : the wake is steady and symmetric
- $Re_w > 44$ : the wake is unsteady and non symmetric with periodic vortex shedding: heat transfer is enhanced
- $Re_w > 140$ : the wake becomes disorderd, heat transfer is further enhanced

#### Frequency response

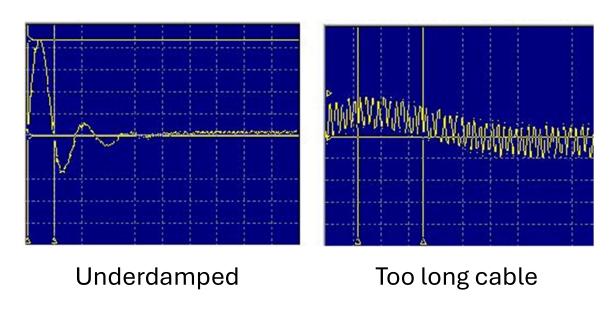


Square excitation

 $\tau$ : empirical response time with undershoot (15%)

 $f_c$ : cut-off frequency

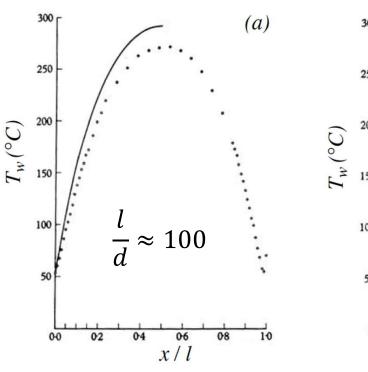
Freymuth P (1977) Frequency response and electronic testing for constant-temperature hot-wire anemometers. J Phys E Sci Instrum 10(7):705–710

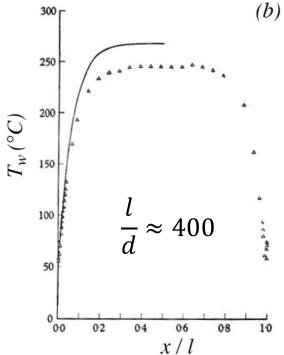


#### Finite length effect

From Champagne et al (1967) reproduced by Lomas (1986)

Non uniform temperature in wire



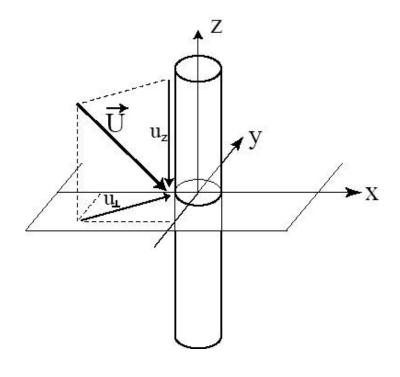


Minimise  $k_w$ Maximise  $\frac{l}{d}$ Keep l not too long

**——** 

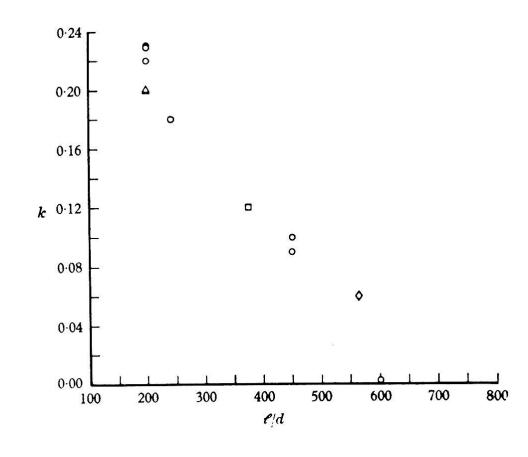
Usual compromise:  $l \approx 1$  mm,  $d \approx 5$   $\mu m$ 

### **Directional sensitivity**



Measured velocity:

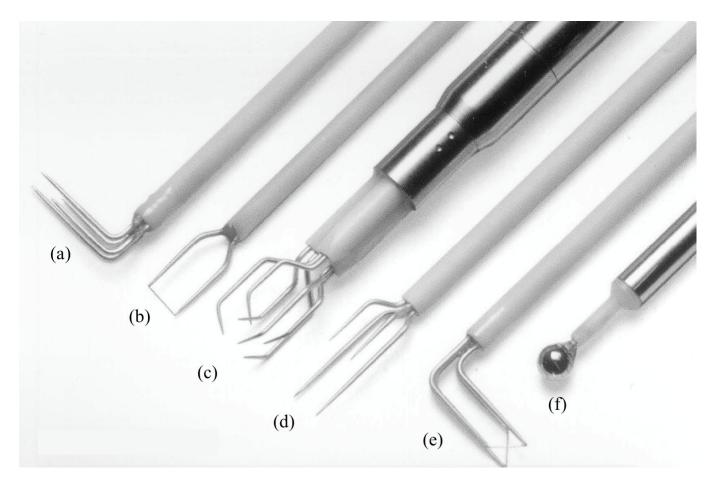
$$U_{eff}^2 = \sqrt{u_\perp^2 + k^2 u_z^2}$$



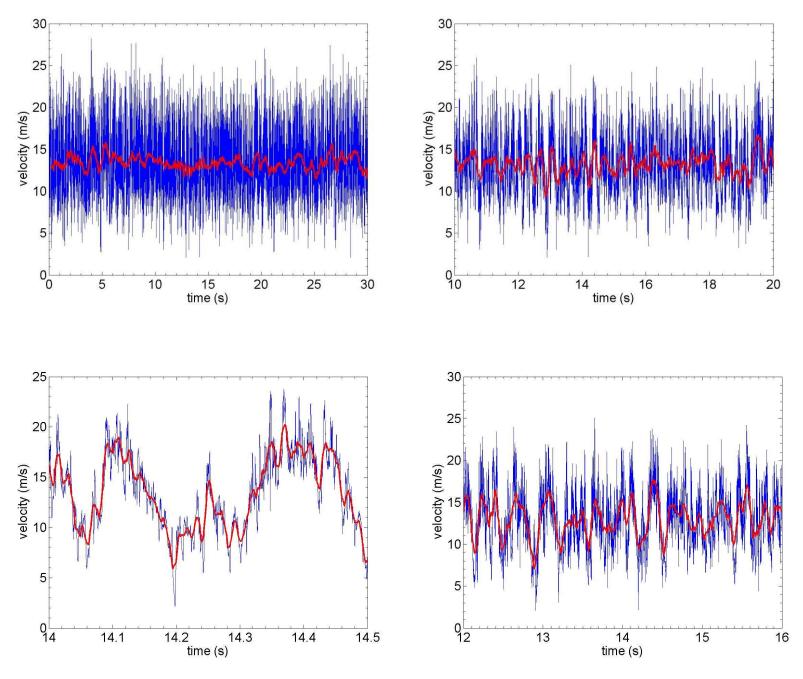
When 
$$\frac{l}{d} \to \infty : k \to 0$$

When 
$$\frac{l}{d} \to 1: k \to 1$$

## **Different probes**



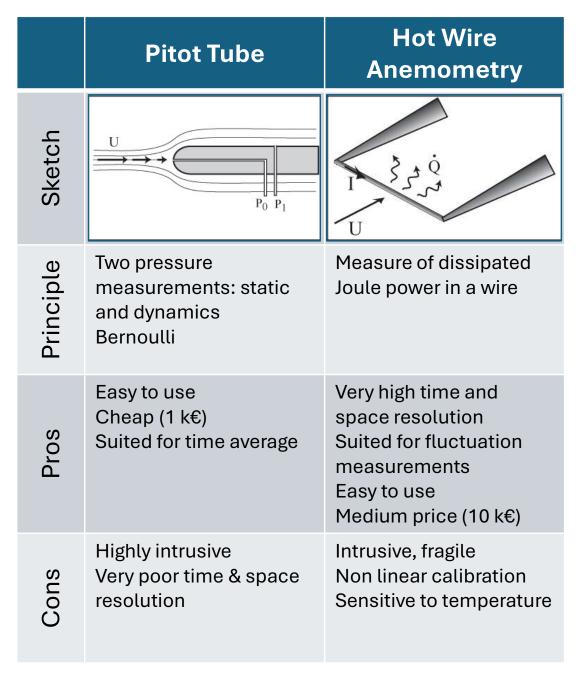
- Several probes (a,c,e): to measure several components
- With cold probes (d): to compensate temperature
- Hot films (f): water measurements



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#### Main features

- Time resolved
- Point measurement (0D)
- Mainstream flow
- Non linear calibration
- Directional ambiguity
- Emperature sensitivity





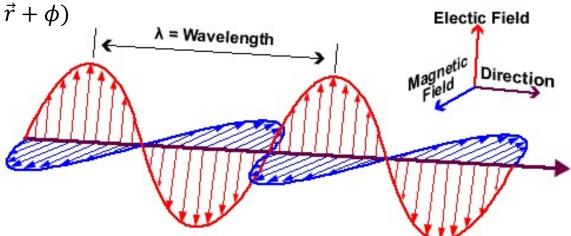
Invented in 1964 by Yeh & Cummins

#### Laser

• Electromagnetic waves

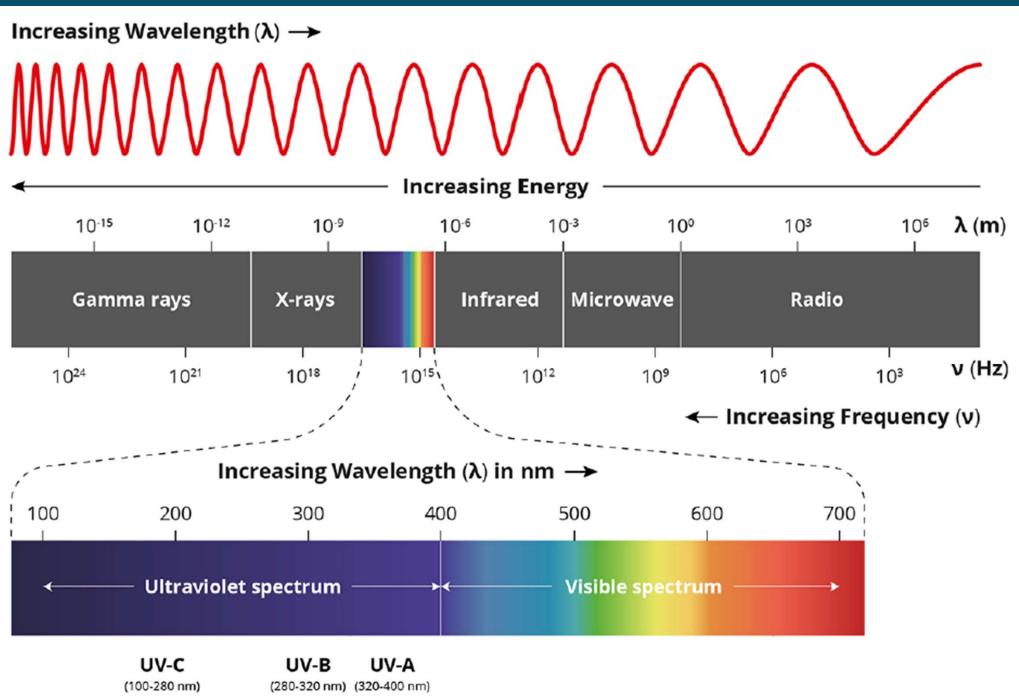
 $\vec{E}(t,\vec{r}) = \vec{p} E_0 \cos(\omega t - \vec{k}.\vec{r} + \phi)$ 

- Wavelength  $\lambda = \frac{2\pi}{k}$
- Celerity  $c = \frac{\omega}{k}$



## **Properties**

- Monochromatic
- Coherent
- Linearly polarized
- Collimated
- Gaussian intensity distribution



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#### **Usual Lasers**

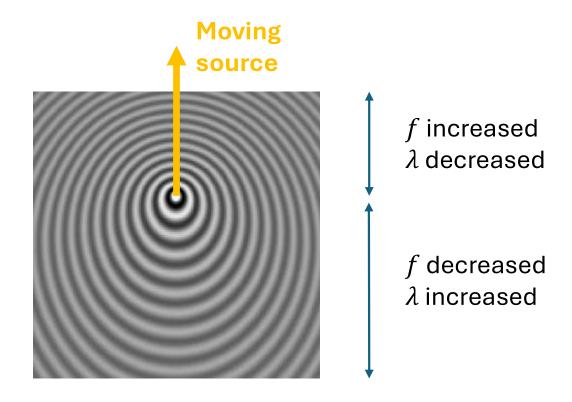
Argon: continuous power in green-blue

Helium: continuous power in red

Nd:Yag: pulsed in infrared shifted to green

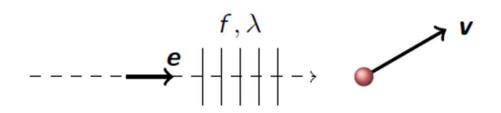
LASER	λ	color	power	diameter
	(nm)		(mW)	(mm)
He-Ne (gas)	632.8	red	1-15	0.65
	476.5	violet	1-600	1.5
Ar <sup>2+</sup> (gas)	488	blue	1-1500	1.5
	514.5	green	1-2000	1.5
doubled YAG (solid)	532	green	20-2000	1

## **Doppler effect**



## Doppler effect

#### Still source, moving receiver

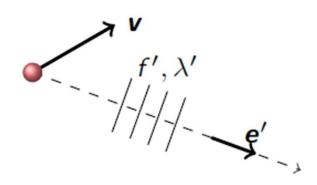


$$T = 1/f$$

$$\lambda' = cT + \vec{v} \cdot \vec{e}T$$

$$f' = f \frac{1}{1 + \frac{\vec{v}}{c} \cdot \vec{e}}$$

#### Moving source, still receiver



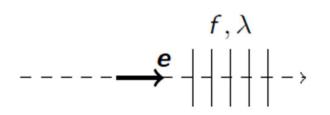
$$T' = 1/f'$$

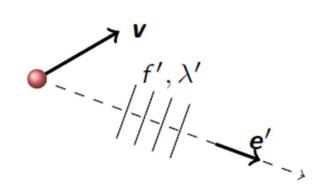
$$\lambda'' = cT' + \vec{v} \cdot \overrightarrow{e'}T'$$

$$f'' = f' \frac{1}{1 - \frac{\vec{v}}{c} \cdot \overrightarrow{e'}}$$

## Doppler effect

#### **Application to velocimetry**





$$f' = f \frac{1}{1 + \frac{\vec{v}}{c} \cdot \vec{e}}$$

$$f'' = f' \frac{1}{1 - \frac{\vec{v}}{c} \cdot \overrightarrow{e'}}$$

#### **Doppler frequency**

$$f^{\prime\prime}\simeq f+f_D$$

$$f_D = f \frac{\vec{v}}{c} \cdot (\overrightarrow{e'} - \overrightarrow{e})$$

#### **Measurement**

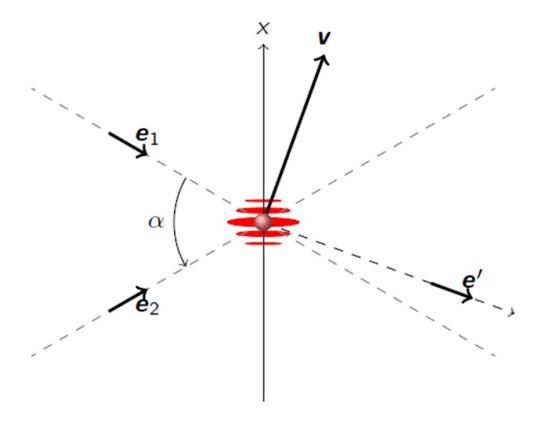
$$f \simeq 10^{14} \ \mathrm{Hz}$$
  $f_D \simeq 10^6 - 10^7 \ \mathrm{Hz}$ 

- Requires a  $10^{-8}$  resolution for 10% accuracy
- Depends on observation direction



Measurement of  $f_D$  is preferred through interferometric system

#### **Doppler effect on crossed beams**



#### Frequency along e':

$$f_1 \simeq f + f \frac{\vec{v}}{c} \cdot (\overrightarrow{e'} - \overrightarrow{e_1})$$

$$f_2 \simeq f + f \frac{\vec{v}}{c} \cdot (\overrightarrow{e'} - \overrightarrow{e_2})$$

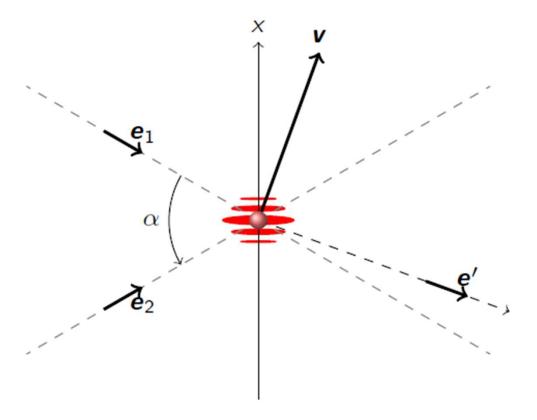
#### Intensity along e':

$$I = E_1^2 + E_2^2 + 2E_1E_2\cos(2\pi f \frac{\vec{v}}{c}.(\vec{e_1} - \vec{e_2}))$$

#### **Measurement**

- One component detected:  $f_D = f \frac{\vec{v}}{c} \cdot (\vec{e_1} \vec{e_2}) = 2f \frac{v_x}{c} \sin \left(\frac{\alpha}{2}\right)$
- No Doppler shift if  $\vec{v} \perp \overrightarrow{e_x}$
- Independent of observation direction

#### **Doppler effect on crossed beams**



#### Measurement

- When  $V_{\chi} = 0$ ,  $f_D = \Delta f$
- When  $V_x < 0$ ,  $f_D < \Delta f$
- When  $V_x > 0$ ,  $f_D > \Delta f$

#### **Problem**

- Insensitive to direction
- Zero velocity not measurable

#### **Solution**

- Scrolling fringes
- Shift frequency  $\Delta f$

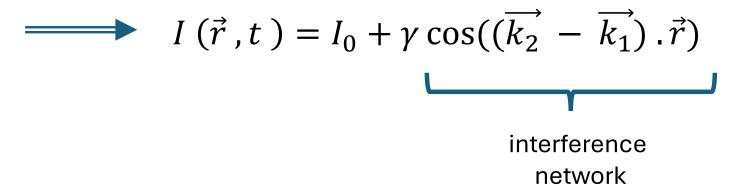
$$f_D = 2f \frac{v_x}{c} \sin\left(\frac{\alpha}{2}\right) + \Delta f$$

#### Interferometry

- Amplitudes sum up
- Intensity is measured

$$I(\vec{r},t) = \left| \vec{E_1} \vec{p_1} \cos(\omega_1 t - \vec{k_1} \cdot \vec{r}) + \vec{E_2} \vec{p_2} \cos(\omega_2 t - \vec{k_2} \cdot \vec{r}) \right|^2$$

#### At the sensor time scale

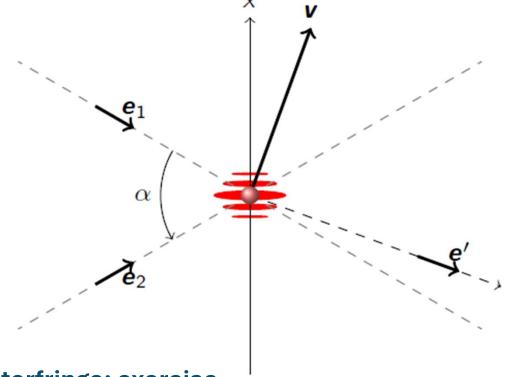


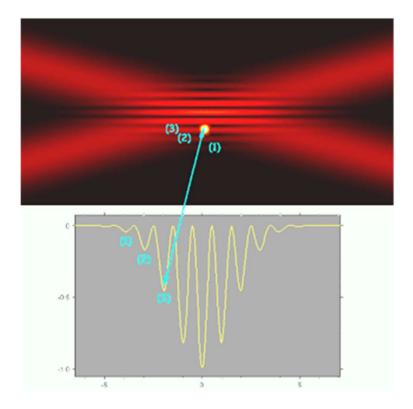
#### Interference pattern

- Spatially structured
- Independent of time
- Max when both have the same polarisation, i.e.  $\overrightarrow{p_1} = \pm \overrightarrow{p_2}$

#### Interferometry

$$I(\vec{r},t) = I_0 + \gamma \cos((\vec{k_2} - \vec{k_1}).\vec{r})$$





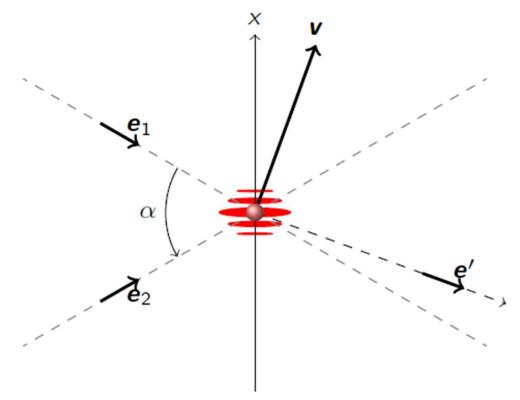
### Interfringe: exercise

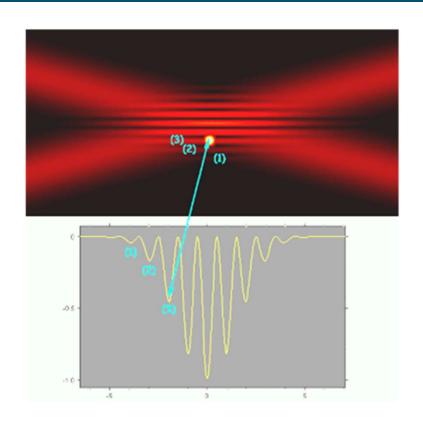
$$(\overrightarrow{k_2} - \overrightarrow{k_1}) \cdot \overrightarrow{r} = 2\pi$$

$$d = \frac{\lambda}{2\sin\frac{\alpha}{2}}$$

#### Interferometry

Seeding particles pass through the fringe at velocity  $ec{v}$ 

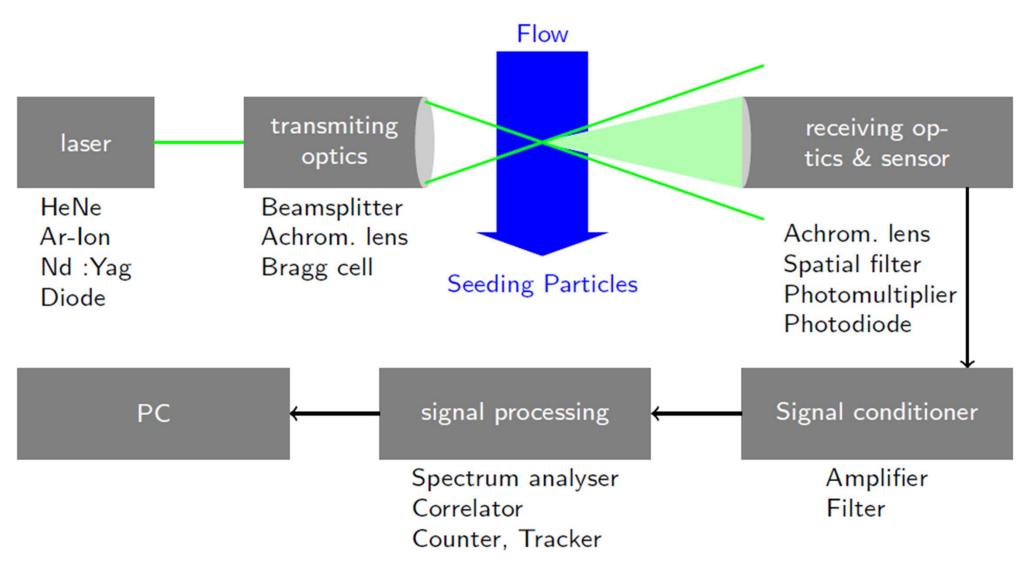




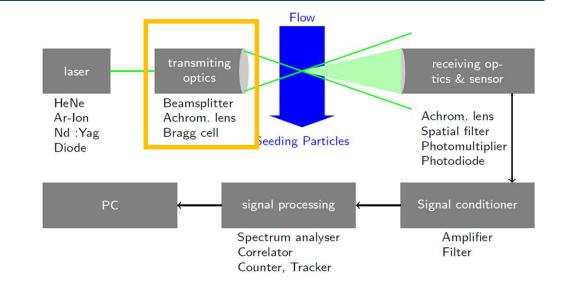
Scattered light is modulated in time at period  $T_D$  or frequency  $f_D$ 

$$T_D = \frac{d}{\overrightarrow{v}.\overrightarrow{e_x}}$$
 We recover the Doppler shift!  $f_D = \frac{2v_x}{\lambda}\sin\frac{\alpha}{2}$  Master 2 – Fluid Mechanics - IPParis

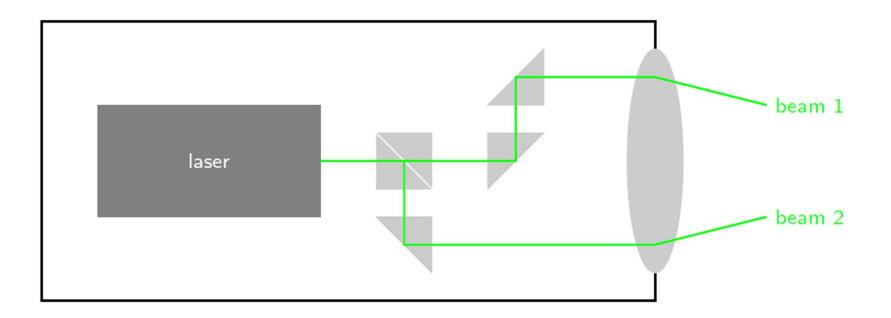
#### **Implementation**



#### **Beam generator**

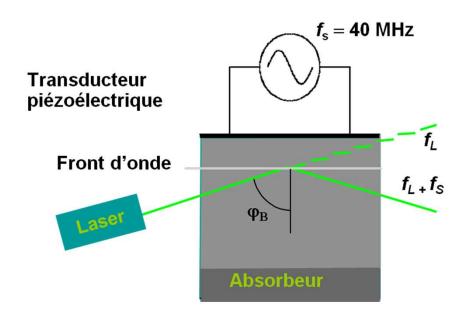


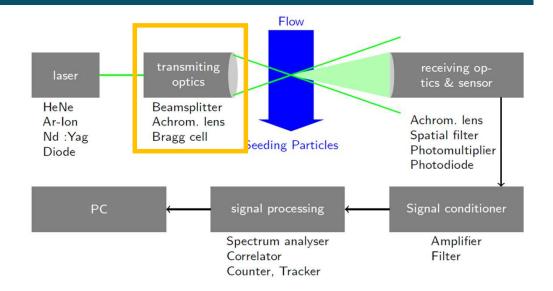
#### Generating two interfering beams

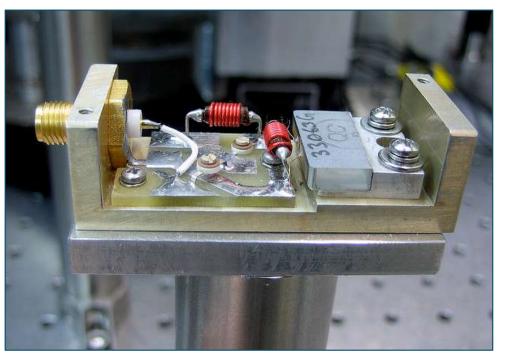


#### **Beam generator**

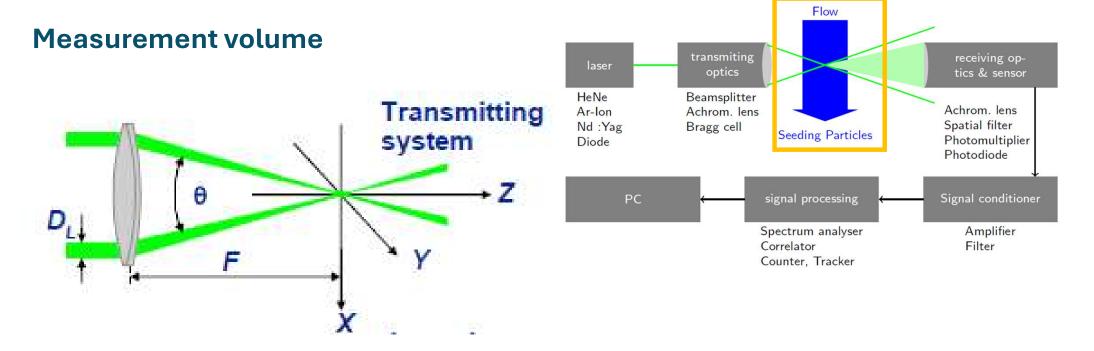
#### **Bragg cell**



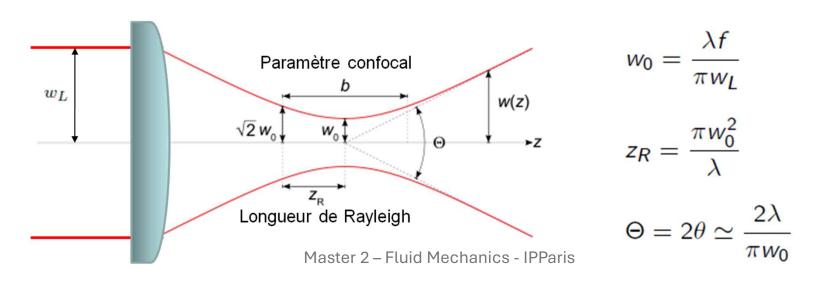




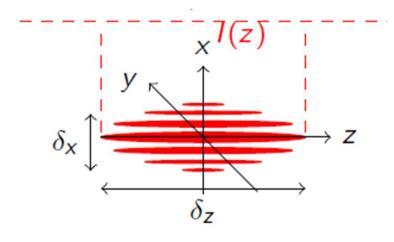
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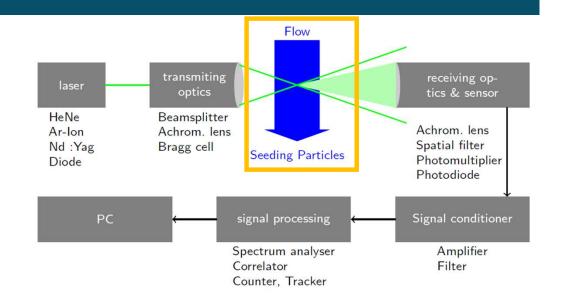


#### Diffraction limit for gaussian beams



#### Measurement volume





Volume dimensions

$$\delta_{z} = \frac{4f\lambda}{\pi D_{L}\sin\frac{\alpha}{2}}, \qquad \delta_{y} = \frac{4f\lambda}{\pi D_{L}}, \qquad \delta_{x} = \frac{4f\lambda}{\pi D_{L}\cos\frac{\alpha}{2}}$$

$$\delta_{y} = \frac{4f\lambda}{\pi D_{L}}$$

$$\delta_X = \frac{4f \,\lambda}{\pi D_L \cos \frac{\alpha}{2}}$$

Interfringe

$$d = \frac{\lambda}{2\sin\frac{\alpha}{2}}$$

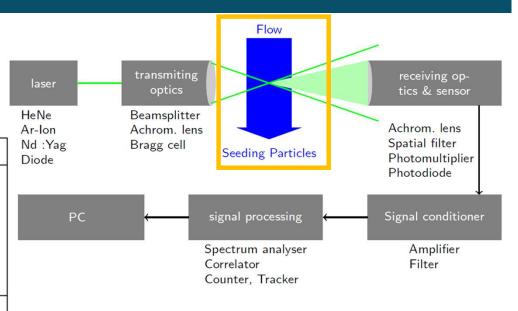
Number of fringes

$$N = rac{8f an rac{lpha}{2}}{\pi D_L}$$

#### **Particles**

#### For liquids

State	Material	Mean diameter $(\mu m)$	
Solid polystyrène		0.5-10	
	Aluminium	2-7	
	Magnesium	2-5	
	synthetic granules	1-10	
	glass microbeads	30-100	
Liquid oils		0.5-10	
	dioctylphathalate	< 1	



#### For gas

State	Material	Mean diameter $(\mu m)$
Solid	polystyrene	10-100
	Aluminium	2-7
	hollow glass sphere	10-100
	granules for synthetic coating	10-500
Liquid	oils	50-500
Gas	bubles of $O_2$ , $H_2$ , etc	50-1000

#### **Particles**

#### Seeking tracers

$$\underbrace{\frac{\pi}{6}\phi^{3}\rho_{p}\frac{d_{p}}{dt}\mathbf{v}_{p}}_{\text{inertial force}} = \underbrace{-3\pi\mu\phi(\mathbf{v}_{p}-\mathbf{v}_{f})}_{\text{Stokes force}} - \frac{\pi\phi^{3}}{6}\underbrace{\nabla P}_{\rho\frac{d\mathbf{v}_{f}}{dt}}$$

$$1\pi_{43} \left(d\mathbf{v}_{f}\right)$$

$$- \frac{1}{2} \frac{\pi}{6} \phi^3 \rho_f \left( \frac{d \mathbf{v}_f}{dt} - \frac{d_p \mathbf{v}_p}{dt} \right)$$

fluid resistance to sphere acceleration

$$+\frac{3}{2}\phi^2\sqrt{\pi\rho_f\mu}\int_0^t\frac{1}{\sqrt{t-t'}}\left(\frac{d\mathbf{v}_f}{dt}-\frac{d_p\mathbf{v}_p}{dt}\right)dt'+\mathbf{f}_{ext}$$

HeNe

Ar-lon

Diode

Nd:Yag

PC

drag force due to an unsteady flow

Maxey, Riley, Gattignol, 1983

receiving op-

tics & sensor

Achrom. lens

Photomultiplier

Signal conditioner

**Amplifier** 

Filter

Spatial filter

Photodiode

Flow

Seeding Particles

signal processing

Spectrum analyser

Counter, Tracker

Correlator

transmiting

optics

Beamsplitter

Achrom. lens

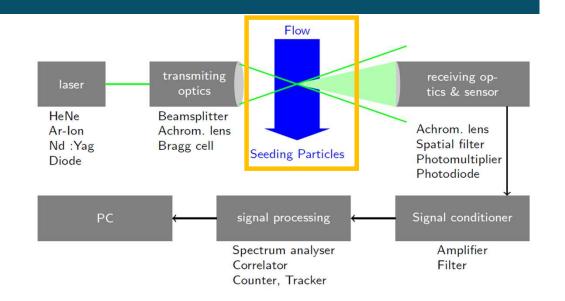
Bragg cell

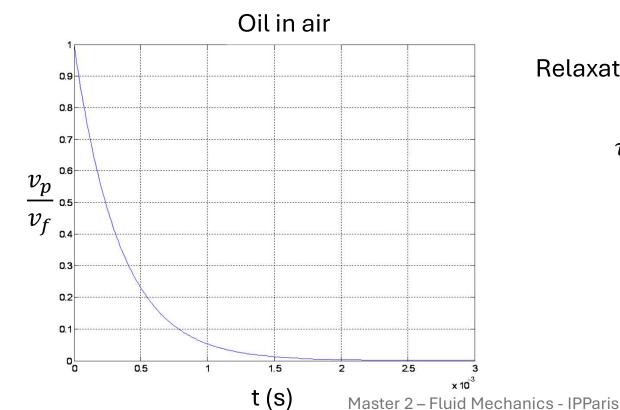
#### **Particles**

Seeking tracers

$$\frac{dv_p}{dt} = 18 \frac{\mu}{\phi^2 \rho_p} \left( v_f - v_p \right)$$

$$v_p = v_f \left( 1 - e^{-t/\tau_p} \right)$$



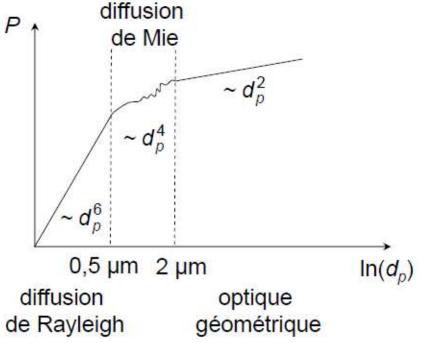


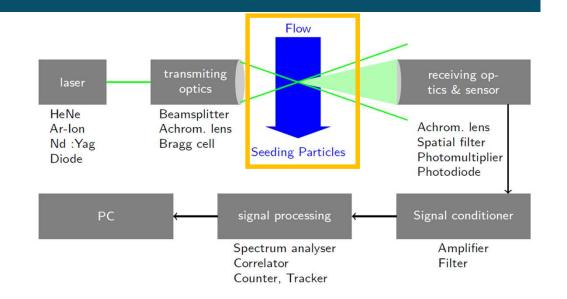
Relaxation time

$$\tau_p = \frac{\rho_p \phi^2}{18\mu}$$

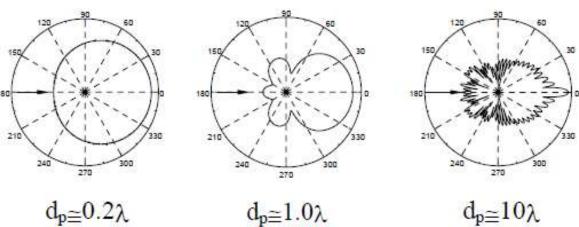
#### **Particles**

Looking for light



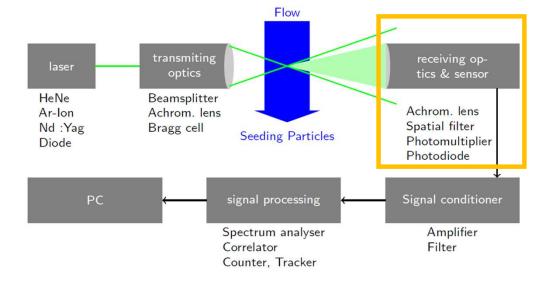


#### Diffracted light in all directions

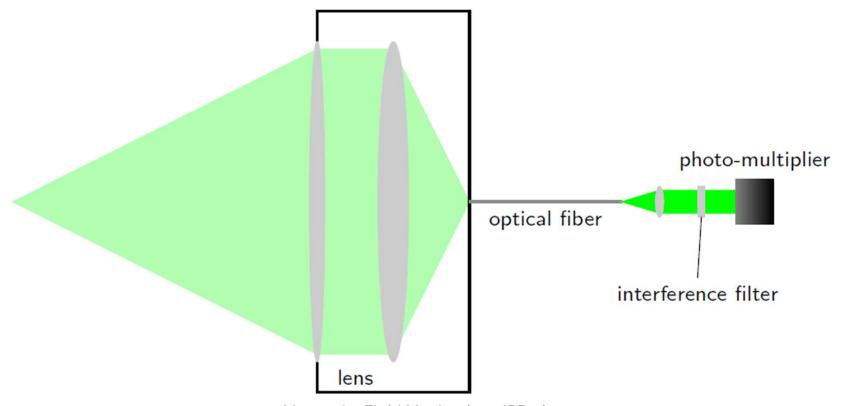


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#### **Implementation**

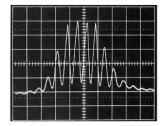


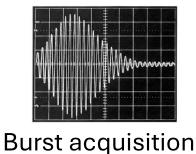
**Detection system** 

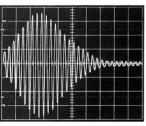


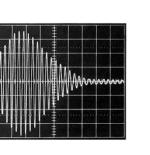
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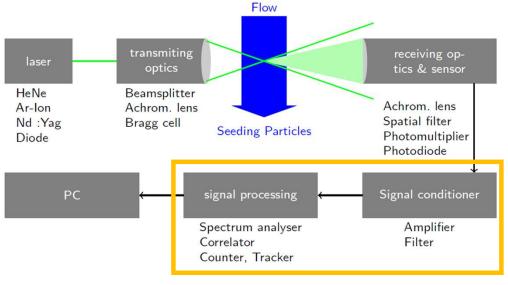
#### **Implementation**



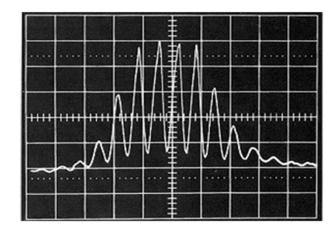






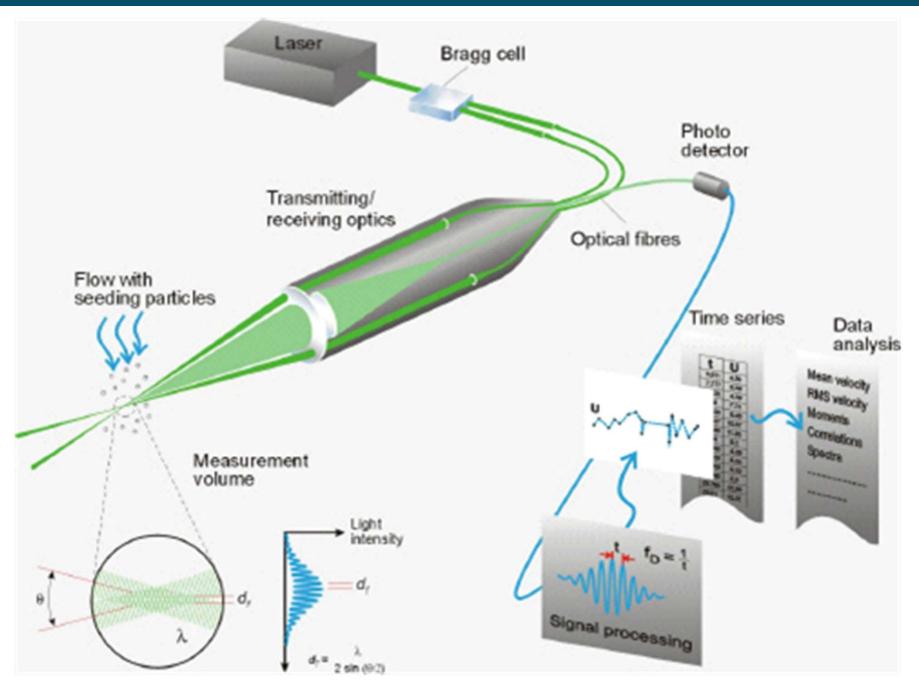


#### Fourier analysis

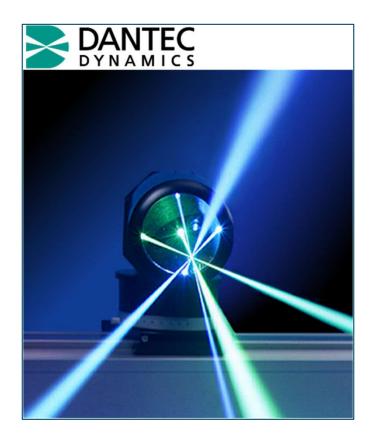


$$I(t) = a \left(1 + \sin(2\pi f_D t)\right) G(t)$$

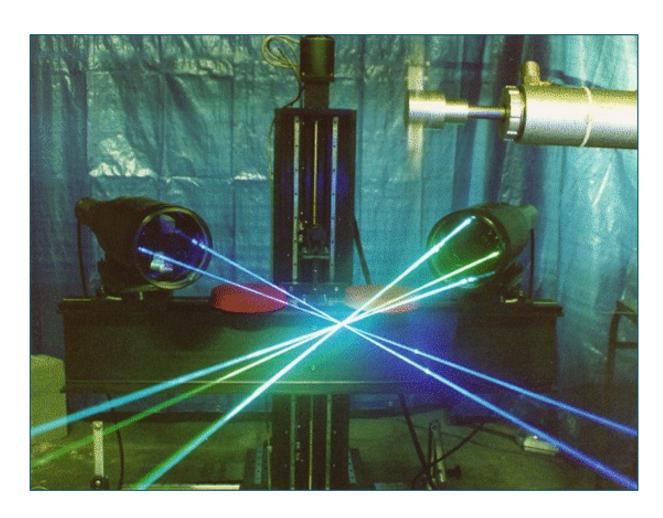
$$|\hat{I}(f)| = a \,\hat{G}(f) + a \,\delta(f - f_D) * \hat{G}(f)$$



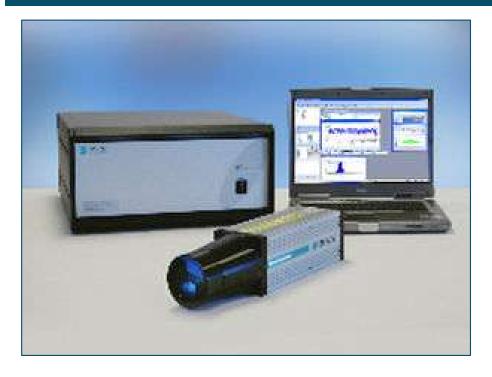
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2D system



3D system



Standard full system



Small integrated 3D FiberFlow probe



Fiber Flow probes (60 & 85 mm)

#### Sources of noise

- Detection noise
- Electronic and thermal noises of the pre-amplifying line.
- High order laser modes (optical noise)
- Diffused light out of the control volume, dirts, damaged window, ambient light,
- Multiple particles, etc.
- Stray reflections (windows, lens, mirrors, ...)



Laser power selection, seeding, optical parameters, ... to optimise the signal over noise ratio

#### **Main features**

- Time-resolved point measurement
- Point measurement is a small ellipsoid (spatial coarsening)
- Non-uniform sampling frequency
- The flow must be seeded
- Linear law
- Expensive

# Comparison

	Pitot Tube	Hot Wire Anemometry	Laser Doppler Anemometry
Sketch	$\begin{array}{c c} U \\ \hline \\ P_0 & P_1 \end{array}$	I Zzzó U	
Principle	Two pressure measurements: static and dynamics Bernoulli	Measure of dissipated Joule power in a wire	Interferometric measurement of a Doppler shift on scattering particle
Pros	Easy to use Cheap (1 k€) Suited for time average	Very high time and space resolution Suited for fluctuation measurements Easy to use Medium price (10 k€)	Non intrusive High time and space resolution Suited for fluctuations Suited for several components
Cons	Highly intrusive Very poor time & space resolution	Intrusive, fragile Non linear calibration Sensitive to temperature	Non regular sampling High price (50-100 k€) Seeding required Difficult settings