

# Solving Maxwell's equations with the Weighted Regularization Method and a Lagrange multiplier

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# Time-dependent Maxwell equations

- In vacuum, over the time interval  $]0, T[$ ,  $T > 0$ .
- Goal: compute the EM field in a domain  $\Omega$  (with Lipschitz polyhedral boundary) encased in a perfect conductor.

Find  $(\mathcal{E}(t), \mathcal{H}(t))$  such that

$$\left\{ \begin{array}{ll} \varepsilon_0 \partial_t \mathcal{E} - \mathbf{curl} \mathcal{H} = -\mathcal{J} & \text{in } \Omega, 0 < t < T ; \\ \mu_0 \partial_t \mathcal{H} + \mathbf{curl} \mathcal{E} = 0 & \text{in } \Omega, 0 < t < T ; \\ \operatorname{div} (\varepsilon_0 \mathcal{E}) = \rho & \text{in } \Omega, 0 < t < T ; \\ \operatorname{div} (\mu_0 \mathcal{H}) = 0 & \text{in } \Omega, 0 < t < T ; \\ \mathcal{E} \times \mathbf{n} = 0 & \text{on } \partial\Omega, 0 < t < T ; \\ \mathcal{E}(0) = \mathcal{E}_0, \mathcal{H}(0) = \mathcal{H}_0 & \text{in } \Omega. \end{array} \right.$$

( Charge conservation equation:  $\partial_t \rho + \operatorname{div} \mathcal{J} = 0$ .

Initial conditions:  $\operatorname{div} \mathcal{E}_0 = \frac{1}{\varepsilon_0} \rho(0)$ ;  $\operatorname{div} \mathcal{H}_0 = 0$ .

$\mathbf{n}$  is the unit outward normal to  $\partial\Omega$ . )

# Related systems of equations

Second order (in time) wave equations...

● In the electric field  $\mathcal{E}$

Equivalent system : Find  $\mathcal{E}(t)$  such that

$$\left\{ \begin{array}{ll} \partial_{tt}^2 \mathcal{E} + c^2 \mathbf{curl} \mathbf{curl} \mathcal{E} = -\frac{1}{\varepsilon_0} \partial_t \mathcal{J} & \text{in } \Omega, 0 < t < T ; \\ \operatorname{div} (\varepsilon_0 \mathcal{E}) = \rho & \text{in } \Omega, 0 < t < T ; \\ \mathcal{E} \times \mathbf{n} = 0 & \text{on } \partial\Omega, 0 < t < T ; \\ \mathcal{E}(0) = \mathcal{E}_0, \partial_t \mathcal{E}(0) = \mathcal{E}_1 & \text{in } \Omega. \end{array} \right.$$

$$\left( \mathcal{E}_1 := \frac{1}{\varepsilon_0} \left( \mathbf{curl} \mathcal{H}_0 - \mathcal{J}(0) \right) \right)$$

Or ...

# Related systems of equations

*Eigenmode computations in a resonator cavity...*

- Assume the time-dependence writes  $\exp(-i\omega t)$ .  
( $\omega > 0$  is the pulsation.)
- *In the electric field  $\mathcal{E}$*

Equivalent system: Find  $(\mathcal{E}, \omega)$  such that

$$\begin{cases} c^2 \mathbf{curl} \mathbf{curl} \mathcal{E} = \omega^2 \mathcal{E} & \text{in } \Omega ; \\ \operatorname{div} \mathcal{E} = 0 & \text{in } \Omega ; \\ \mathcal{E} \times \mathbf{n} = 0 & \text{on } \partial\Omega. \end{cases}$$

Or ...

# Related systems of equations

*(Magnetic) quasi-static computations...*

- Assume that the electric displacement current  $\varepsilon_0 \partial_t \mathcal{E}$  is negligible.
- *In the electric field  $\mathcal{E}$*

Find  $\mathcal{E}$  such that

$$\left\{ \begin{array}{ll} \mathbf{curl} \mathcal{E} = -\mu_0 \partial_t \mathcal{H} & \text{in } \Omega, 0 < t < T ; \\ \operatorname{div} (\varepsilon_0 \mathcal{E}) = \rho & \text{in } \Omega, 0 < t < T ; \\ \mathcal{E} \times \mathbf{n} = 0 & \text{on } \partial\Omega, 0 < t < T. \end{array} \right.$$

# Functional space: possible choices

- Which *functional space* to measure the electric field?

First choice:

$$\mathcal{H}_0(\mathbf{curl}, \Omega) := \{\mathcal{F} \in L^2(\Omega)^3 \mid \mathbf{curl} \mathcal{F} \in L^2(\Omega)^3, \mathcal{F} \times \mathbf{n}|_{\partial\Omega} = 0\}.$$

(cf. [Kikuchi'87/'89], [Demkowicz et al'9x], [Boffi et al'9x/'0x], ...)

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(cf. [Kikuchi'87/'89], [Demkowicz et al'9x], [Boffi et al'9x/'0x], ...)

Scalar product:  $(u, v)_{\mathcal{H}(\mathbf{curl}, \Omega)} := (u, v)_0 + (\mathbf{curl} u, \mathbf{curl} v)_0.$

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$$\mathcal{X}_0 := \{\mathcal{F} \in \mathcal{H}_0(\mathbf{curl}, \Omega) \mid \operatorname{div} \mathcal{F} \in L^2(\Omega)\}.$$

OK in a convex domain  $\Omega$

(cf. [Assous-Degond-Heintz -Raviart-Segr '93].)

OK in a 2D or 2D1/2 non-convex domain  $\Omega$  (**Singular Complement Method**)

(cf. [Assous-Jr et al'98/'00/'03], [Bonnet-Hazard-Lohrengel'99/'02].)

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# Functional space: possible choices

- Which *functional space* to measure the electric field?

Third choice:

$$\mathcal{X}_\gamma := \{\mathcal{F} \in \mathcal{H}_0(\mathbf{curl}, \Omega) \mid \operatorname{div} \mathcal{F} \in L^2_\gamma(\Omega)\}.$$

$$\left( L^2_\gamma(\Omega) := \{v \in L^2_{\text{loc}}(\Omega) \mid w_\gamma v \in L^2(\Omega)\}, \|v\|_{0,\gamma} := \|w_\gamma v\|_0. \right.$$

The **weight**  $w_\gamma$  is a function of the distance  $r$  to the reentrant edges (called  $E$ ):

$$w_\gamma(r) \approx r^\gamma \text{ for small } r,$$

with a suitable  $\gamma \in ]\gamma_{\min}, 1[$ ,  $0 < \gamma_{\min} < \frac{1}{2}$ , cf. [\[Costabel-Dauge'02/'03\]](#). )

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This is the so-called **Weighted Regularization Method**: *our choice from now on...*

# The constraint on the divergence

- What happens if one wants to take into account the *constraint* on the divergence of the electric field *explicitly*?  $\left( \operatorname{div}(\varepsilon_0 \mathcal{E}) = \rho \text{ or } \operatorname{div} \mathcal{E} = 0. \right)$

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Motivations:

- Improve the quality of the divergence of the discrete fields.  
 $\left( \text{For instance, for the computed eigenmodes.} \right)$
- Resolve numerical problems related to the discrete charge conservation equation.  
 $\left( \text{Solve the Vlasov-Maxwell system to compute the motion of charged particles.} \right)$

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Solution:

Introduce a *Lagrange multiplier*.

# Mixed Variational Formulations

- The *eigenproblem* to be solved writes equivalently ( $\lambda = \omega^2/c^2$ )  
Find  $(\mathcal{E}, \lambda) \in \mathcal{K}_\gamma \times \mathbb{R}^+$  such that

$$(\mathbf{curl} \mathcal{E}, \mathbf{curl} \mathcal{F})_0 = \lambda(\mathcal{E}, \mathcal{F})_0, \quad \forall \mathcal{F} \in \mathcal{K}_\gamma,$$

with  $\mathcal{K}_\gamma := \{\mathcal{F} \in \mathcal{X}_\gamma \mid \operatorname{div} \mathcal{F} = 0\}$ .

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Find  $(\mathcal{E}, \lambda) \in \mathcal{K}_\gamma \times \mathbb{R}^+$  such that

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# Mixed Variational Formulations

- The *mixed eigenproblem* to be solved writes  
Find  $(\mathcal{E}, p, \lambda) \in \mathcal{X}_\gamma \times L^2_{-\gamma}(\Omega) \times \mathbb{R}^+$  such that

$$\begin{cases} (\mathcal{E}, \mathcal{F})_{\mathcal{X}_\gamma} + L^2_{-\gamma} \langle p, \operatorname{div} \mathcal{F} \rangle_{L^2_\gamma} = \lambda(\mathcal{E}, \mathcal{F})_0, \quad \forall \mathcal{F} \in \mathcal{X}_\gamma \\ L^2_{-\gamma} \langle q, \operatorname{div} \mathcal{E} \rangle_{L^2_\gamma} = 0, \quad \forall q \in L^2_{-\gamma}(\Omega). \end{cases}$$

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- A *discrete approximation* is  $((\mathcal{X}_h)_h \subset \mathcal{X}_\gamma, (M_h)_h \subset L^2_{-\gamma}(\Omega))$   
Find  $(\mathcal{E}_h, p_h, \lambda_h) \in \mathcal{X}_h \times M_h \times \mathbb{R}^+$  such that

$$\begin{cases} (\mathcal{E}_h, \mathcal{F}_h)_{\mathcal{X}_\gamma} + L^2_{-\gamma} \langle p_h, \operatorname{div} \mathcal{F}_h \rangle_{L^2_\gamma} = \lambda_h (\mathcal{E}_h, \mathcal{F}_h)_0, \quad \forall \mathcal{F}_h \in \mathcal{X}_h \\ L^2_{-\gamma} \langle q_h, \operatorname{div} \mathcal{E}_h \rangle_{L^2_\gamma} = 0, \quad \forall q_h \in M_h. \end{cases}$$

Abstract convergence theory, see [\[Boffi-Brezzi-Gastaldi'97\]](#), [\[Boffi'06\]](#).

Uses *strong approximability* of solutions  $\mathcal{E}$ , *weak approximability* of solutions  $p$   
(with  $(\mathcal{E}, p)$  solutions to the *plain* mixed problem...)

# Elements of convergence theory

- A desired property is the **uniform discrete inf-sup condition**

$$\exists \beta > 0, \forall h, \inf_{q_h \in M_h} \sup_{\mathcal{F}_h \in \mathcal{X}_h} \frac{L_{-\gamma}^2 \langle q_h, \operatorname{div} \mathcal{F}_h \rangle_{L_{\gamma}^2}}{\|\mathcal{F}_h\|_{\mathcal{X}_{\gamma}} \|q_h\|_{0, -\gamma}} \geq \beta.$$

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- Formulation with  $\mathcal{E}$  set in  $\mathcal{X}_0$  ( $\Omega$  convex or SCM in 2D, 2D1/2 domains).
  - With the  $P_2 - iso - P_1$  Taylor-Hood finite element, as in [Assous-Degond-Heintz -Raviart-Segr '93]. The *udisc* is satisfied, cf. [Girault-Jr'02].
  - With the  $P_{k+1} - P_k$  Taylor-Hood finite elements, the *udisc* is satisfied, cf. [Stenberg'84], [Boffi'97].

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- Formulation in  $\mathcal{E}$  set in  $\mathcal{X}_{\gamma}$  (WRM).
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The *udisc* is not satisfied anymore!

Why?

# A negative result

- In order to check the discrete inf-sup condition, let

$$\beta_h = \inf_{q_h \in M_h} \sup_{\mathcal{F}_h \in \mathcal{X}_h} \frac{L_{-\gamma}^2 \langle q_h, \operatorname{div} \mathcal{F}_h \rangle_{L_\gamma^2}}{\|\mathcal{F}_h\|_{\mathcal{X}_\gamma} \|q_h\|_{0, -\gamma}}.$$

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How can one estimate  $(\beta_h)_h$ ?

- Introduce the *plain* mixed Variational Formulation (**rhs**  $f, g$ ).  
Find  $(\mathcal{E}_h, p_h)$  such that

$$\begin{cases} a(\mathcal{E}_h, \mathcal{F}_h) + b(p_h, \mathcal{F}_h) = f(\mathcal{F}_h), \quad \forall \mathcal{F}_h \in \mathcal{X}_h \\ b(q_h, \mathcal{E}_h) = g(q_h), \quad \forall q_h \in M_h. \end{cases}$$

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How can one estimate  $(\beta_h)_h$ ?

- Introduce the *matrix version* of the plain mixed Variational Formulation.  
Find  $(\vec{\mathcal{E}}, \vec{p})$  such that

$$\begin{cases} \mathbb{A} \vec{\mathcal{E}} + \mathbb{B}^T \vec{p} = \vec{f} \\ \mathbb{B} \vec{\mathcal{E}} = \vec{g}. \end{cases}$$

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- Proposition (e. g. [Jamelot'05](#)): Define  $\mathbb{M}$  by  $(\mathbb{M}\vec{q} | \vec{q}) = \|q_h\|_{M_h}^2$ . There holds

$$\kappa(\mathbb{M}^{-1}(\mathbb{B} \mathbb{A}^{-1} \mathbb{B}^T)) \leq \left( \frac{\|b\|}{\beta_h} \right)^2.$$

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How can one estimate  $(\beta_h)_h$ ?

- Practical experiments with the  $P_2 - P_1$  Taylor-Hood finite element.
  - In the unit cube (see [\[Hechme-Jr'07a\]](#))

Meshsize	$h'$	$h'/2$	$h'/4$
$\kappa$	2.9	2.8	2.8

⇒ *Consistent* with the fact that  $(\beta_h)_h$  is independent of  $h$ ...

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How can one estimate  $(\beta_h)_h$ ?

- Practical experiments with the  $P_2 - P_1$  Taylor-Hood finite element.
  - The WRM in a 2D L-shape domain (see [\[Hechme-Jr'07a\]](#))

Meshsize	$h$	$h/2$	$h/4$	$h/8$
$\kappa$	29	69	161	364

⇒  $(\beta_h)_h$  decreases sharply when  $h$  decreases...

# A new family of finite elements for the WRM

- Consider a family of triangular/tetrahedral meshes  $(\mathcal{T}_h)_h$  of  $\Omega \subset \mathbb{R}^d$ .

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$$\mathcal{X}_h = \{\mathcal{F}_h \in C^0(\bar{\Omega})^d \mid \mathcal{F}_h|_T \in P_{k+1}(T)^d, \forall T \in \mathcal{T}_h, \text{ and } \mathcal{F}_h \times \mathbf{n}|_{\partial\Omega} = 0\},$$

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- *New family of  $P_{k+1} - P_k$  finite elements (cf. [\[Hechme-Jr'07a\]](#)):*

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with  $E_h$  a neighborhood of the reentrant corners and/or edges:

$$E_h = \cup_{T \in \mathcal{T}_h \text{ s.t. } T \cap E \neq \emptyset} T.$$

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$\implies$

Zero Near Singularity  $P_{k+1} - P_k$  finite elements.

# Remarks

- In the variational formulations, at the discrete level, one has

$$L^2_{-\gamma} \langle \bar{q}_h, \operatorname{div} \mathcal{F}_h \rangle_{L^2_\gamma} = (\bar{q}_h, \operatorname{div} \mathcal{F}_h)_0, \forall (\mathcal{F}_h, q_h) \in \mathcal{X}_h \times \bar{M}_h \dots$$

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As a consequence, the quantity of interest is

$$\bar{\beta}_h = \inf_{\bar{q}_h \in \bar{M}_h} \sup_{\mathcal{F}_h \in \mathcal{X}_h} \frac{(\bar{q}_h, \operatorname{div} \mathcal{F}_h)_0}{\|\mathcal{F}_h\|_{\mathcal{X}_\gamma} \|\bar{q}_h\|_{0, -\gamma}}.$$

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  - Follows (more or less!) the series of lemmas of [Stenberg'84].
  - Difficulties:
    - Presence of *weights* in  $\|\mathcal{F}_h\|_{\mathcal{X}_\gamma}$  and  $\|\bar{q}_h\|_{0, -\gamma}$ .
    - *Local estimates* (near the reentrant edges).
    - Existence of *gradients* in  $\mathcal{X}_h$ .
    - Non-zero mean value Lagrange multipliers.

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- The final result on  $(\bar{\beta}_h)_h$ :  
Measuring the quality of the *regular* family of triangulations  $(\mathcal{T}_h)_h$ .

$$\exists \sigma > 1, \forall h, \forall T \in \mathcal{T}_h, h_T \leq \sigma \rho_T ;$$

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$\implies$  Zero Near Singularity finite elements satisfy the udisc.

# Computing eigenvalues and eigenvectors

- Find  $(\mathcal{E}_h, \bar{p}_h, \lambda_h) \in \mathcal{X}_h \times \bar{M}_h \times \mathbb{R}^+$  such that

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  - $(E_\lambda)_{\lambda \leq \lambda_n}$  the corresponding *eigenspaces*.
  - Approximation error  $\varepsilon_\lambda(h) = \sup_{v \in E_\lambda, \|v\|_{\mathcal{X}_\gamma} = 1} \inf_{\mathcal{F}_h \in \mathcal{X}_h} \|v - \mathcal{F}_h\|_{\mathcal{X}_\gamma}$   
(worst case:  $\varepsilon_\lambda(h) \leq C_\varepsilon h^{\gamma - \gamma_{\min} - \varepsilon}$ .)
  - Error on eigenvalues:  $|\lambda - \lambda_h| < C_n \varepsilon_\lambda(h)^2$ .
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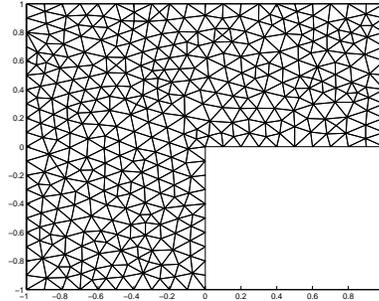
$\implies$  No more spurious eigenmodes.

# Numerical experiments in 2D

On a 'practical' example, taken from [Monique Dauge's](#) benchmark.

- 2D, L-shaped, domain, straight sides, corners in  $(0,0)$ ,  $(1,0)$ ,  $(1,1)$ ,  $(-1,1)$ ,  $(-1,-1)$ ,  $(0,-1)$ .
- First five eigenvalues (with repetition), up to six digits:
  - $\lambda_1 = 1.47562$ , eigenmode has the **strong unbounded singularity**;
  - $\lambda_2 = 3.53403$ ;  $\lambda_3 = \lambda_4 = 9.86960$ ;  $\lambda_5 = 11.3895$ .
- The weight is implemented with  $\gamma = 0.95$  (NB.  $\gamma_{min} = 1/3$ .)
- Experiments (cf. [\[Buffa-Jamelot-Jr'07\]](#)):
  - on a series of *quasi-uniform meshes*;
  - relative errors  $r_{k,h} = |\lambda_{k,h} - \lambda_k|/\lambda_k$ ,  $1 \leq k \leq 5$  are reported.

# Numerical experiments in 2D



- Three meshes with respectively
  - 738, 2952 and 11808 triangles ;
  - 410, 1557 and 6065 vertices ;
- Results for the Zero Near Singularity finite elements:

<i>mesh</i>	$r_{1,h}$	$r_{2,h}$	$r_{3,h}$	$r_{4,h}$	$r_{5,h}$
<i>uniform1</i>	$1.3e - 2$	$3.3e - 4$	$9.4e - 5$	$1.1e - 4$	$9.9e - 3$
<i>uniform2</i>	$8.0e - 3$	$6.2e - 5$	$2.3e - 5$	$2.5e - 5$	$1.3e - 5$
<i>uniform3</i>	$4.4e - 3$	$1.2e - 5$	$5.5e - 6$	$6.2e - 6$	$5.3e - 6$

# Comparisons in 3D

On a second 'practical' example, taken from [Monique Dauge's](#) benchmark.

- 3D, thick L-shaped, domain  $(] - 1, 1[^2 \setminus [-1, 0]^2) \times ]0, 1[$ .
- First nine eigenvalues (with repetition), up to six digits:
  - $\lambda_1 = 9.6397$ ;  $\lambda_2 = 11.3452$ ;  $\lambda_3 = 13.4036$ ;  $\lambda_4 = 15.1972$ ;
  - $\lambda_5 = 19.5093$ ;  $\lambda_6 = \lambda_7 = \lambda_8 = 19.7392$ ;  $\lambda_9 = 21.2591$ .
- The weight is implemented with  $\gamma = 0.95$  (NB.  $\gamma_{min} = 1/3$ .)
- Experiments (cf. [\[Hechme-Jr'07b\]](#)):
  - on a *graded mesh* (grading towards the reentrant edge);
  - relative errors  $r_{k,h} = |\lambda_{k,h} - \lambda_k|/\lambda_k$ ,  $1 \leq k \leq 9$  are reported;

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    - comparison of the mixed approach with
      - the *parameterized approach* [\[Costabel-Dauge'02\]](#) with parameter  $s = \nu$ .
- Find  $(\mathcal{E}'_h, \lambda'_h) \in \mathcal{X}_h \times \mathbb{C}$  such that

$$(\mathbf{curl} \mathcal{E}'_h, \mathbf{curl} \mathcal{F}_h)_0 + \nu(\operatorname{div} \mathcal{E}'_h, \operatorname{div} \mathcal{F}_h)_{0,\gamma} = \lambda'_h(\mathcal{E}'_h, \mathcal{F}_h)_0, \quad \forall \mathcal{F}_h \in \mathcal{X}_h.$$

*Spurious* (curl-free) eigenvalues are **filtered out** by comparing

$$\operatorname{Re}(\lambda'_h) \text{ to } \operatorname{Im}(\lambda'_h).$$

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  - Experiments (cf. [\[Hechme-Jr'07b\]](#)):
    - comparison of the mixed approach with
      - the *filter approach* [\[Costabel-Dauge'03\]](#), [\[Hechme-Jr'07b\]](#).
- Find  $(\mathcal{E}_h, \lambda_h) \in \mathcal{X}_h \times \mathbb{R}^+$  such that

$$(\mathcal{E}_h, \mathcal{F}_h)_{\mathcal{X}_\gamma} = \lambda_h (\mathcal{E}_h, \mathcal{F}_h)_0, \quad \forall \mathcal{F}_h \in \mathcal{X}_h.$$

*Spurious* (curl-free) eigenvalues are **filtered out** by evaluating the *filter ratio*

$$\frac{\|\operatorname{div} \mathcal{E}_h\|_{0,\gamma}}{\|\mathbf{curl} \mathcal{E}_h\|_0}.$$

# Comparisons in 3D

- A mesh with  
4032 tetrahedra; 1010 vertices.
- Number of d.o.f.  
15818 for the parameterized and filter approaches; 18162 for the mixed approach.
- Results:

Method	Filter	Parameterized	Mixed
$r_1$	$6.1 \times 10^{-4}$	$6.1 \times 10^{-4}$	$6.2 \times 10^{-4}$
$r_2$	$6.5 \times 10^{-3}$	$1.1 \times 10^{-2}$	$8.5 \times 10^{-3}$
$r_3$	$8.1 \times 10^{-4}$	$7.4 \times 10^{-4}$	$8.4 \times 10^{-4}$
$r_4$	$1.1 \times 10^{-4}$	$1.0 \times 10^{-4}$	$1.1 \times 10^{-4}$
$r_5$	$2.0 \times 10^{-3}$	$4.7 \times 10^{-3}$	$6.9 \times 10^{-3}$
$r_6$	$1.8 \times 10^{-4}$	$1.8 \times 10^{-4}$	$1.8 \times 10^{-4}$
$r_7$	$1.2 \times 10^{-3}$	$1.1 \times 10^{-3}$	$1.2 \times 10^{-3}$
$r_8$	$1.2 \times 10^{-3}$	$1.1 \times 10^{-3}$	$1.3 \times 10^{-3}$
$r_9$	$1.3 \times 10^{-3}$	$1.1 \times 10^{-3}$	$1.1 \times 10^{-2}$

# Focusing on eigenvalues or eigenvectors?

On a last 'practical' example, taken from [Monique Dauge's](#) benchmark.

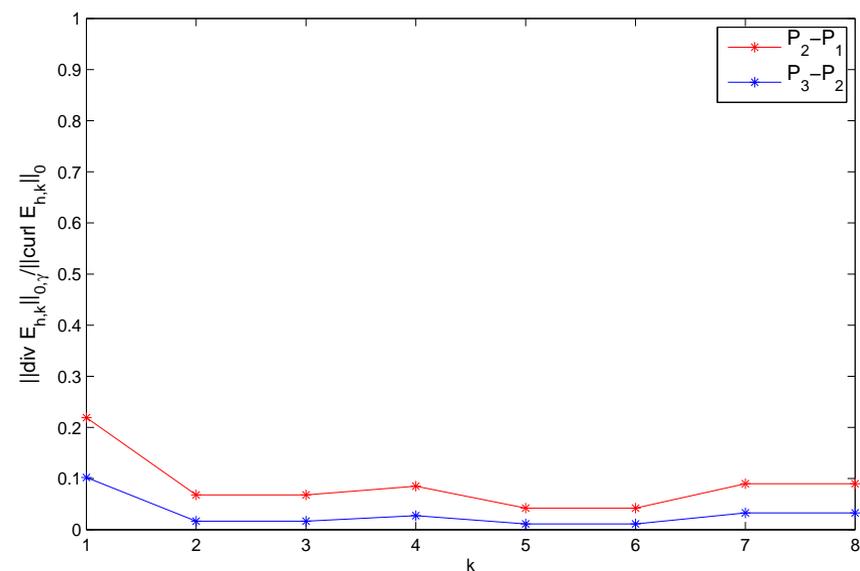
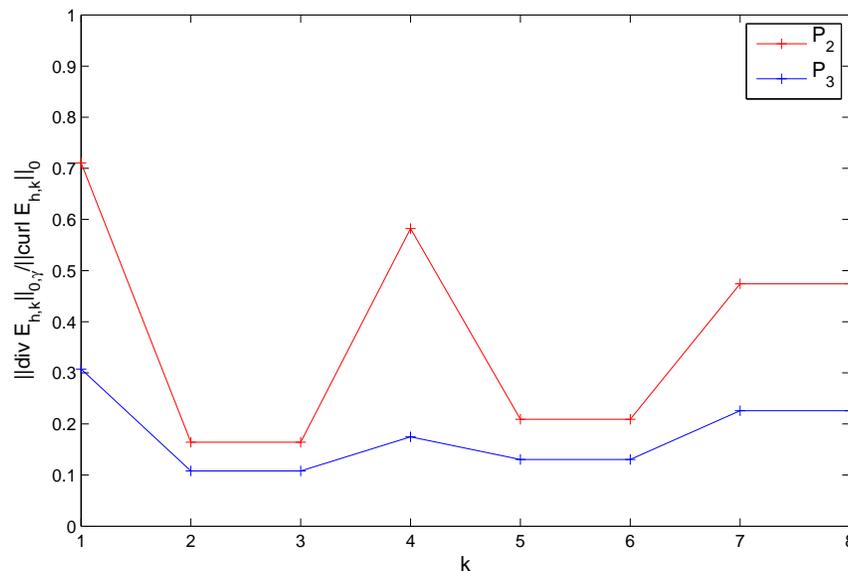
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- The weight is implemented with  $\gamma = 0.95$  (NB.  $\gamma_{min} = 1/3$ .)
- A graded mesh with  
2688 tetrahedra; 665 vertices.
- Experiments on the first eight eigenpairs (cf. [\[Hechme-Jr'07b\]](#)):

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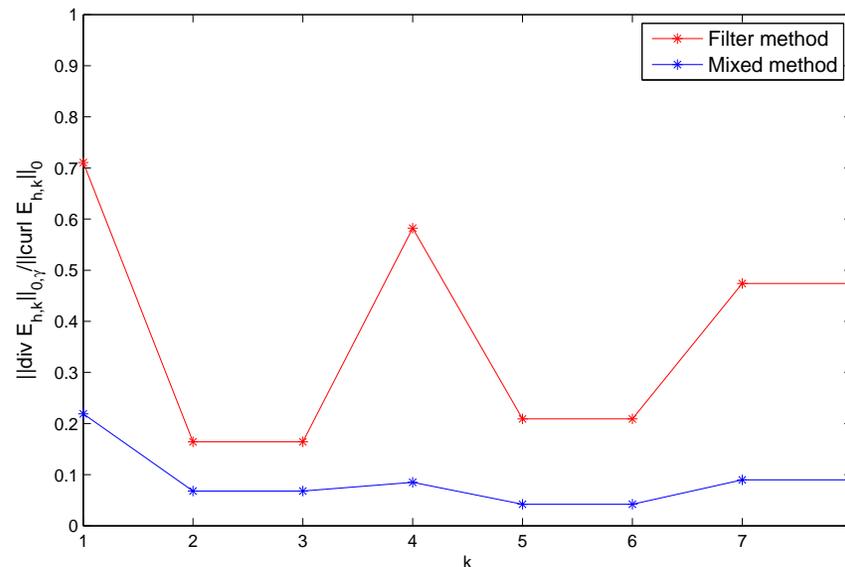
Filter ratios for the filter (left) and mixed (right) methods



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Filter ratios for both methods ( $P_2$  FE for the field)



# Concluding remarks

- Implementing the mixed method with the WRM turned out to be a challenging problem!
- The classical  $P_{k+1} - P_k$  Taylor-Hood finite elements fail to verify the udisc. The Zero Near Singularity  $P_{k+1} - P_k$  finite elements provide an adequate answer. (with [G. Hechme](#).)
- These FE allowed us to solve accurately the EM eigenvalue problem in mixed form. No more spurious eigenmodes. (with [E. Jamelot](#), [A. Buffa](#), [G. Hechme](#).)

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- These FE allowed us to solve accurately the EM eigenvalue problem in mixed form. No more spurious eigenmodes. (with [E. Jamelot](#), [A. Buffa](#), [G. Hechme](#).)
- Application to the time-dependent problem (Vlasov-Maxwell) has been completed. (with [S. Labrunie](#).)
- Extension to materials ( $\varepsilon, \mu$  piecewise constant) is possible. (with [F. Lefèvre](#), [S. Lohrengel](#), [S. Nicaise](#).)